A Physics-Based Parameterization of Air–Sea Momentum Flux at High Wind Speeds and Its Impact on Hurricane Intensity Predictions

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ABSTRACT

A new bulk parameterization of the air–sea momentum flux at high wind speeds is proposed based on coupled wave–wind model simulations for 10 tropical cyclones that occurred in the Atlantic Ocean during 1998–2003. The new parameterization describes how the roughness length increases linearly with wind speed and the neutral drag coefficient tends to level off at high wind speeds. The proposed parameterization is then tested on real hurricanes using the operational Geophysical Fluid Dynamics Laboratory (GFDL) coupled hurricane–ocean prediction model. The impact of the new parameterization on the hurricane prediction is mainly found in increased maximum surface wind speeds, while it does not appreciably affect the hurricane central pressure prediction. This helps to improve the GFDL model–predicted wind–pressure relationship in strong hurricanes. Attempts are made to provide physical explanations as to why the reduced drag coefficient affects surface wind speeds but not the central pressure in hurricanes.

1. Introduction

The transfer of momentum between the atmosphere and the ocean is a crucial subject in atmospheric modeling and weather forecasting. In most global and mesoscale atmospheric models the momentum exchange at the sea surface is parameterized using the drag coefficient \( C_d \) that increases approximately linearly with wind speed. This behavior of the \( C_d \) is based on extrapolations from field measurements in weak-to-moderate wind regimes less than 25 m s\(^{-1}\). In strong wind conditions, it has been reported that the Geophysical Fluid Dynamics Laboratory (GFDL) hurricane prediction model, used operationally by the U.S. National Hurricane Center, tends to underestimate the surface wind speed for a given central pressure (Ginis et al. 2004, see their Fig. 2). This may be partly attributed to an insufficient horizontal grid resolution. However, a more likely reason is the inadequate representation of the physical processes that affect air–sea momentum exchanges at high wind speeds.

Recently, Moon et al. (2004a,b,c) used a coupled wave–wind (CWW) model to show that the drag coefficient levels off (or even decreases) at wind speeds exceeding 30 m s\(^{-1}\). This behavior of the \( C_d \) at high wind speeds was also found in recent field observations of Powell et al. (2003), laboratory experiments of Alamaro et al. (2002) and Donelan et al. (2004), and theoretical studies of Emanuel (2003) and Makin (2005). Although physical explanations vary and are not conclusive, there is a general consensus that the \( C_d \) ceases to increase with wind speed at high wind speeds. Therefore, the present parameterization of the \( C_d \) in atmospheric models, where the \( C_d \) linearly increases with wind speed, clearly overestimates the momentum flux at high wind speeds. The overestimated momentum flux likely contributes to the limited skill of numerical hurricane intensity prediction. Reduced \( C_d \) at high winds will likely have a significant impact not only on hurricane modeling but also on other atmospheric and oceanic modeling focusing on high wind conditions. However, no studies have been performed to investigate the effect of reduced \( C_d \) at high wind speeds on atmospheric and oceanic modeling.

Since direct flux measurements at very high wind speeds are still extremely limited (Kepert 2004; French 2005), theoretical and numerical approaches may pro-
vide a useful and complementary way to develop new flux parameterizations in a wide range of wind conditions without the limitations of available observational data. This study proposes a simple parameterization of air–sea momentum fluxes for high wind conditions, based on the CWW model simulations (Moon et al. 2004c) for 10 tropical cyclones in the Atlantic Ocean during 1998–2003. The new parameterization is compared with available observations and results from other theoretical/numerical studies, and is applied to the GFDL coupled hurricane–ocean prediction system to investigate its effect on real hurricane simulations.

In section 2, the main features of the CWW model will be briefly discussed since the present approach strongly depends on the results of the CWW model simulations. A new parameterization based on hurricane simulation results is proposed in section 3. Section 4 describes the numerical experiments for real hurricanes using the new flux parameterization. The summary and conclusion are given in the last section.

2. CWW model

In the CWW model the surface wave directional frequency spectrum near the spectral peak is calculated using the WAVEWATCH-III (Tolman 2002) model and the high-frequency part of the spectrum is parameterized using the theoretical model of Hara and Belcher (2002). The complete wave spectrum is then introduced to the wave boundary layer model of Hara and Belcher (2004) to estimate the drag coefficient at different wind and wave conditions. The CWW model treats the wind stress as a vector quantity to consider the influence of dominant waves that propagate at a large angle to the local wind. Therefore, the method allows us to estimate the wind stress vectors for any given surface wave field, even for the complex seas encountered under tropical cyclones. The details of the model are provided in Moon et al. (2004a,b).

The CWW model has been applied for constant winds from 10 to 45 m s$^{-1}$ (Moon et al. 2004a), idealized hurricanes with various forward speeds (Moon et al. 2004b), and 10 real hurricanes in the Atlantic Ocean during 1998–2003 (Moon et al. 2004c). The main results of these studies are that the neutral drag coefficient levels off at high wind speeds and the nondimensional roughness length is mainly determined by two parameters: the input wave age (wave age determined by the peak frequency of wind energy input) and the wind speed, regardless of the complexity of the wave field. The model-simulated drag coefficients are generally consistent with the recent results of field observations, laboratory experiments, and theoretical studies at high wind speeds (Powell et al. 2003; Alamaro et al. 2002; Donelan et al. 2004; Emanuel 2003).

3. New parameterization of air–sea momentum flux

At low winds (wind speed less than 12.5 m s$^{-1}$) the bulk parameterization used in the operational GFDL hurricane model is consistent with previous observations. Therefore, the roughness length ($z_0$) in our parameterization is identical to the bulk parameterization and is calculated using a constant Charnock coefficient or nondimensional roughness length ($z_{ch}$), defined as

$$z_{ch} = \frac{z_0 g}{u_{*}^2} = 0.0185,$$

(1)

where $u_*$ is the friction velocity and $g$ is the gravitational acceleration. In neutral conditions the neutral wind speed ($W$) at 10-m height is described by the logarithmic profile,

$$W = \frac{u_*}{\kappa} \ln \frac{10}{z_0},$$

(2)

where $\kappa$ is the von Kármán constant (=0.4) and the unit of $z_0$ is meters. From Eqs. (1) and (2), the friction velocity $u_*$ (m s$^{-1}$) can be expressed as a function of $W$ (m s$^{-1}$) using polynomial fitting,

$$u_* = 0.001 W^2 + 0.028 W, \quad W \leq 12.5 \text{ m s}^{-1},$$

(3)

where the regression coefficient is 0.99. The roughness length can be also expressed as a function of $W$:

$$z_0 = \frac{0.0185}{g} \left(0.001 W^2 + 0.028 W\right), \quad W \leq 12.5 \text{ m s}^{-1}.$$  

(4)

Our parameterization of $z_0$ at higher winds speeds ($W > 12.5$ m s$^{-1}$) is based on the CWW simulations of 10 tropical cyclones that occurred in the Atlantic Ocean during 1998–2003, reported by Moon et al. (2004c). Their scatterplot of calculated $z_0$ against wind speed is reproduced in Fig. 1a (plus signs). We then establish an empirical relationship between $z_0$ and $W$ by fitting the CWW results with a linear regression function:

$$z_0 = (0.085 W - 0.58) \times 10^{-3}, \quad W > 12.5 \text{ m s}^{-1},$$

(5)

where the regression coefficient is 0.87 (see Fig. 1a, blue line).
The expressions (4) and (5) match well at $W = 12.5$ m s$^{-1}$ with the difference of $1 \times 10^{-6}$ and together make up our new parameterization of $z_0$ for all wind speeds. In most ocean models the air–sea momentum flux is calculated using the neutral drag coefficient $C_d$ (Flather 1994; Tang et al. 1996). An expression of the neutral $C_d$ as a function of $W$ can be obtained by introducing (4) or (5) into

$$C_d = \kappa^2 \left( \ln \frac{10}{z_0} \right)^{-2}.$$  

(6)

In many atmospheric and weather prediction models, the air–sea momentum flux is calculated using $z_0$ as a function of $u_*$ (Charnock 1955; Kurihara and Tuleya 1974). Such a formulation is obtained by first expressing $W$ as a function of $u_*$ from (2) and (5),

$$W = -0.56u_*^2 + 20.255u_* + 2.458,$$  

(7)

using polynomial fitting (where the regression coefficient is 0.99), and then combining (1), (5), and (7):

$$z_0 = \frac{0.0185}{g} u_*^2, \quad W \leq 12.5 \text{ m s}^{-1} \quad \text{and} \quad (8a)$$

$$z_0 = [0.085(-0.56u_*^2 + 20.255u_* + 2.458)$$

$$- 0.58] \times 10^{-3}, \quad W > 12.5 \text{ m s}^{-1}.$$  

(8b)

Figure 1b compares the neutral drag coefficient estimated from our new formulation (thick blue line) with the bulk parameterization used in the operational GFDL model, the results from Wu (1982), Large and Pond (1981), Donelan et al. (2004), and Powell et al. (2003). For $W \leq 12.5$ m s$^{-1}$, the new $C_d$ represents a monotonic increase with wind speed as in the operational GFDL model and is similar to that of Wu (1982). But it is slightly higher than those of Large and Pond (1981) and Donelan et al. (2004). For $W > 12.5$ m s$^{-1}$, the new $C_d$ tends to level off between 2 and 3. This is similar to the trend observed by Donelan et al. (2004) and is within the error bars estimated by Powell et al. (2003), although it is somewhat higher than their averaged values (squares in Fig. 1b). At 60 m s$^{-1}$ wind speed, the new $C_d$ is half of the value used in the operational GFDL model.

It should be noted that at high wind speeds the boundary layer close to the surface is nearly neutral since $\ln z/z_0 \gg f_m(z/L)$, where $f_m$ is the universal function and $L$ is the Monin–Obukhov (MO) length. At low winds, however, the stability effect may be important. In the following numerical experiments the parameterization in (8) is used to define the roughness length, and then the real 10-m wind speed and exchange coefficients are calculated through the iteration using stability parameters and universal functions based on the MO similarity theory (Liu et al. 1979; DeCosmo et al. 1996).
4. Impact of the new momentum flux parameterization on the GFDL model hurricane predictions

The proposed parameterization has been tested for real hurricane predictions using the operational GFDL hurricane prediction model, which is coupled with the Princeton ocean model. The GFDL model is a primitive equation model formulated in latitude, longitude, and sigma coordinates with 42 levels in the vertical. Major features of the model are given by Kurihara et al. (1998). The model includes a multiply nested movable mesh configuration and model initialization by the method of vortex replacement (Kurihara et al. 1993). The 2005 version of the GFDL model, which has a movable innermost mesh of $1/12^\circ$, is used in the present experiments. Five-day forecasts are conducted for 11 cases of Hurricanes Isabel (2003), Ivan (2004), Frances (2004), Jeanne (2004), and Charley (2004), that is, four cases for Isabel (0000 UTC 10 September, 1800 UTC 18 September, 0000 UTC 12 September, and 0000 UTC 12 September), four cases for Ivan (0000 UTC 10 September, 0600 UTC 10 September, 0000 UTC 11 September, and 0000 UTC 12 September), and one case each for Frances (0600 UTC 1 September), Jeanne (0000 UTC 19 September), and Charley (1800 UTC 11 August). The initial forecast times of each hurricane are carefully selected to consider various situations. During the forecast periods, Hurricanes Isabel and Ivan reached up to category 5 on the Saffir–Simpson scale. Charley (category 4) made landfall and underwent both developing and decaying stages. Frances was in a weakening stage from category 4 to 1, while Jeanne was in a developing stage from category 1 to 2.

Two forecast runs are conducted for each case. In both experiments, the flux calculations are made using the GFDL model original bulk parameterization based on the MO similarity theory. The neutral drag coefficient is expressed by Eq. (6). The neutral exchange coefficient of heat and water (assumed equal, $C_k$) is expressed as (DeCosmo et al. 1996; Zeng et al. 1998; Kurihara and Tuleya 1974)

$$C_k = k^2 \left( \ln \frac{10}{z_0} \right)^{-1} \left( \ln \frac{10}{z_T} \right)^{-1},$$

where $z_T$ (m) is the roughness length for heat and humidity fluxes (assumed equal).

In the first control run (operational GFDL model), both $z_0$ and $z_T$ are defined by Charnock’s relation, 0.0185$u^2_*$, of the formulation in (1). In the second run, $z_T$ is defined by the same Charnock’s relation, but $z_0$ is defined by new empirical formulation in (8). This implies that $C_k$ in the second run is also decreased due to the reduced values of $z_0$ by the new formula at high wind speeds (see Fig. 2), although the heat and humidity parameterizations are not modified here. It is certainly possible that the value of $z_T$ at high wind speeds is also different from the Charnock’s relation. However, the objective of this study is to focus on the effect of decreased $z_0$, and therefore $z_T$ is kept unchanged.

The ratio of the neutral drag and heat/humidity coefficients, $C_k/C_d$, is a key parameter in tropical cyclone intensity (Emanuel 1995). Figure 3 shows the scatterplot of $C_k/C_d$ for two runs. In the first run, $C_k/C_d$ is 1, while in the second run it lies in the range 0.7–1.4 and increases as wind speed increases. The latter trend is qualitatively consistent with Emanuel (1995) who shows, using a hurricane model, that the ratio lies in the range 0.75–1.5 and increases with maximum wind when hurricanes intensify. This, however, contradicts the studies by Liu et al. (1979) and Geernaert et al. (1987), who show that the ratio $C_k/C_d$ is less than 1 and decreases with increasing wind. As Bao et al. (2002) point out, the behavior of $C_k/C_d$ at high winds is still controversial due to lack of reliable observation data under high wind conditions.

The effect of the new flux parameterization on hurricane predictions is investigated in terms of the forecast skill for the surface maximum wind speed (MWS),
central pressure, and track of a hurricane. Figure 4 shows an example of 5-day forecast run for Hurricane Ivan (initial forecast time: 0000 UTC 12 September 2004) using the two flux parameterizations. Here, Hurricane Ivan reached category 5 and sustained its strength for a few days before weakening. The major differences between the two experiments are found in the MWS forecast. The run using the new formula predicts overall higher wind speeds (about 10% increase) than the operational model with the largest difference of about 10 m s\(^{-1}\). The results with the new formula are more consistent with the observations, therefore indicating the improved forecast skill. However, such distinct difference between the two experiments is not found in the central pressure prediction (Fig. 4b). The track forecast is changed only slightly by using the new formula. To investigate how the new formula affects the structure of the hurricane, the surface wind fields at 1800 UTC 12 September from the two models are compared in Fig. 5 with the real-time surface wind analysis of the National Oceanic and Atmospheric Administration’s (NOAA’s) Hurricane Research Division (HRD; Powell et al. 1998) at 1930 UTC 12 September. The figure shows that the asymmetric wind structure appearing in the HRD wind is well reproduced in both model simulations, but that the wind structure near the eyewall, particularly the position and shape of the strong wind area, is better simulated by the new formula.

Another example of the improvements in the MWS forecasts with the new flux parameterization is presented in Fig. 6 for Hurricane Frances (initial forecast time: 0000 UTC 1 September 2004). In this case, Hurricane Frances was in a weakening stage from category 4 to 1 and making landfall at Florida. The overall results for Frances are very similar to the case of Hurricane Ivan, that is, the maximum wind speed increases but the
central pressure changes little with the new flux formula. While it is expected that the maximum wind is increased by reducing the surface drag, it is less clear why the central pressure is not appreciably changed. The pressure–wind relationship is one of the most difficult parameters to predict and it is likely sensitive to the model surface layer physics and the convection parameterization.

Figure 7 compares the SST as well as the surface momentum flux (MF), heat/humidity flux (HF), and energy dissipation rate, averaged within 100-km radius of the hurricane center for the Ivan case shown in Fig. 4, in the two models. Figure 7a shows that the maximum SST drop reached 2.5°C when Ivan entered the Gulf of Mexico and again before landfall. This is consistent with Walker et al. (2005) who observed a significant surface cooling (3°C–7°C) in two large areas along Ivan’s track in the Gulf of Mexico. Considering that the SST drop in Fig. 7a is averaged over a 100-km radius around the hurricane center, it appears that the model-simulated SST cooling is reasonable. It is seen that the MF is significantly reduced by the new formula (Fig. 7b) and this seems to cause an increase of the MWS. On the other hand, the averaged HF (Fig. 7c; i.e., energy input) is not changed much because it is balanced by both the positive effects (from less cooling due to the reduced MF and from increased wind speeds) and the negative effect (from the decrease of the heat transfer coefficient). This implies that the SST cooling plays a significant role in this balance, which is an interesting and important feature of the fully coupled atmosphere–ocean system.

For a mature hurricane the amount of total kinetic energy generated is equal to that being dissipated by friction. The dissipation rate per unit area is the MF times the wind speed (Emanuel 1999). A comparison of downward kinetic energy fluxes (energy dissipation) between the two models (Fig. 7d) shows that the change of energy due to the new parameterization is relatively small (the reduction of the MF is compensated by the increase of the wind speed), particularly after 44 h. This may explain the similar central pressures between the two models in Figs. 4b and 6b.
For all 11 cases, predictions of the central pressure and the MWS between 2 models are compared in Fig. 8. The figure shows that the new flux parameterization tends to increase the MWS but does not affect the central pressure under a wide range of conditions. Figure 9 compares the MWS from 2 models with observations during 11 forecasts of 5 hurricanes. This shows that the predicted MWS with the new formula are more consistent with observations than the one from the operational model. In Fig. 10, a pressure–wind relationship derived from observations is compared with the predictions from 2 models for the same 11 cases. It is clearly seen that the underestimation of the MWS in the operational GFDL model is significantly improved by using the new flux parameterization, in particular, for strong hurricanes.

Therefore, it is expected that the new flux parameterization contributes to the overall improvement of the hurricane MWS forecast in the GFDL model, although more real case forecasts are necessary to confirm these results.
5. Summary and conclusions

In most atmospheric and oceanic models, the air–sea momentum flux at high wind speeds has been parameterized using a bulk formula based on extrapolation from field measurements in weak-to-moderate wind regimes less than 25 m s$^{-1}$. The bulk formula yields a monotonic increase of the drag coefficient ($C_d$) with wind speed. However, recent observational, laboratory, theoretical, and modeling studies for high wind speeds (Powell et al. 2003; Donelan et al. 2004; Emanuel 2003; Moon et al. 2004a,b,c) suggest that $C_d$ levels off (or even decreases) in high wind conditions.

This study proposes a new and rather simple air–sea momentum flux parameterization, which is consistent with the recent results at high wind speeds. The new parameterization is derived from the CWW model simulations of 10 tropical cyclones in the Atlantic Ocean during 1998–2003 (Moon et al. 2004c). The
The new momentum flux parameterization has been tested in 11 forecasts of 5 real hurricanes in the Atlantic Ocean using the operational GFDL hurricane model. While in these experiments the heat and humidity flux calculations are made using the GFDL model original bulk parameterization, the heat/humidity coefficients ($C_h$) are decreased due to the reduced values of $z_0$ by the new formula at high wind speeds. The results show that the new parameterization mainly contributes to an increase of the maximum wind speed (MWS) predictions due to the reduction of the momentum fluxes, but does not appreciably affect the central pressure predictions in these simulations. The analysis of momentum/heat fluxes and energy dissipation as well as sea surface temperature cooling near the hurricane center reveals that the heat flux remains virtually unchanged by the new momentum parameterization because it is balanced by both the positive effects (from less sea surface cooling due to the reduced momentum flux and from increased wind speeds) and the negative effect (from the decrease of the heat transfer coefficient). It is an important feature of the fully coupled atmosphere–ocean system that the hurricane-induced SST change plays a significant role in this balance.

The numerical experiments for five hurricanes show that the new parameterization contributes to overall improvement of hurricane MWS predictions as well as pressure–wind relationships prediction, especially for strong hurricanes. In this study, however, the experiments were conducted only for the GFDL hurricane model which suffers from the underestimation of the MWS for the given central pressure. The underestimation of the MWS in high winds may not be true for other numerical prediction models. More simulations with various models will be necessary to substantiate the results reported here.

It is believed that the air–sea momentum flux is not a function of wind speed alone, but depends on the sea state (Toba et al. 1990; Smith et al. 1992). Under hurricanes the sea state varies significantly according to the relative position from the storm center even under the same wind speed (Moon et al. 2003, 2004b). Therefore, the consideration of sea state can lead to an asymmetric distribution for exchange coefficients and fluxes, which may play a significant role in determining the wind structure inside the hurricane core region and may affect hurricane forecasts. However, the consideration of

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